

QUENCHING FOR MULTI-DIMENSIONAL SEMILINEAR PARABOLIC PROBLEMS ON A BALL WITH A LOCALIZED SOURCE

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ABSTRACT. We study the quenching set of a multi-dimensional semilinear parabolic problem on a ball subject to the first initial-boundary condition. The source term of this problem is a nonlinear localized function. This function tends to infinity when the solution u approaches a finite number. This mathematical model illustrates a nonlinear reaction of a dynamical system occurring at a single location. The main result of this paper is that u quenches at a single point only and the blow-up set of u_t is the whole domain.

Keywords. Localized Source; Quenching; Green's Function

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1. INTRODUCTION

Let $T \in (0, \infty]$, a be a positive real number, c be a positive constant, x_0 be a fixed point in \mathbb{R}^n where $n = 1, 2, \dots$, and $B_1(x_0)$ be an n -dimensional open ball such that $B_1(x_0) = \{x \in \mathbb{R}^n : \|x - x_0\| < 1\}$ centered at x_0 with radius 1, where $\|x - x_0\|$ represents Euclidean distance between the points x and x_0 . We denote the closure and boundary of $B_1(x_0)$ by $\overline{B_1(x_0)}$ and $\partial B_1(x_0)$, respectively. We would like to study the first initial-boundary value problem for the following multi-dimensional semilinear parabolic differential equation with a localized source at x_0 :

$$(1.1) \quad u_t - \Delta u = af(u(x_0, t)) \text{ in } B_1(x_0) \times (0, T),$$

$$(1.2) \quad u(x, 0) = 0 \text{ for } x \in \overline{B_1(x_0)}, u(x, t) = 0 \text{ for } x \in \partial B_1(x_0) \text{ and } t > 0.$$

A solution u is said to quench at (x^*, T) if there exists a sequence (x_n, t_n) such that $x_n \rightarrow x^*$, $t_n \rightarrow T^-$, and $u(x_n, t_n) \rightarrow c^-$ as $n \rightarrow \infty$. Here, x^* is the quenching point, and T is the quenching time.

Physically, equation (1.1) describes instabilities in a system with a localized reaction (cf. [1, 8]). Without the localized source, the forcing function with $a = 1$ is

$f(u(x, t))$. When $f(u) = 1/(1-u)$, the concept of quenching in one spatial dimension was introduced by Kawarada [6] in 1975.

Throughout this paper, we assume that $f(u)$ satisfies the hypothesis below:

(H) $f(u) \in C^2([0, c))$, $f > 0$, $f' > 0$, $f'' > 0$, and $f(u) \rightarrow \infty$ as $u \rightarrow c^-$.

The purpose of this paper is to prove that u_t blows up for all $x \in B_1(x_0)$ and u quenches only at the single point x_0 in some finite time T . In Section 2, we prove some properties of u . In Section 3, we prove that u quenches in some finite time when a is sufficiently large. Based on this result, we prove that u_t tends to infinity in the whole domain and u quenches at x_0 only.

2. PROPERTIES OF THE SOLUTION

Let L be the parabolic operator given by $Lu = u_t - \Delta u$, and $b(x, t)$ be a nontrivial and nonnegative bounded function on $\overline{B_1(x_0)} \times [0, \infty)$. We prove the following comparison theorem.

Lemma 2.1. *Assume that U is a classical solution of the following problem*

$$LU \geq b(x, t)U(x_0, t) \text{ in } B_1(x_0) \times (0, T)$$

$$U(x, 0) \geq 0 \text{ for } x \in \overline{B_1(x_0)}, U(x, t) \geq 0 \text{ for } x \in \partial B_1(x_0) \text{ and } t > 0,$$

then $U(x, t) \geq 0$ on $\overline{B_1(x_0)} \times [0, T)$.

Proof. Let ε and β be positive real numbers and $W(x, t) = U(x, t) + \varepsilon e^{\beta t}$. Then, $W(x, 0) > 0$ for $x \in \overline{B_1(x_0)}$, $W(x, t) > 0$ for $x \in \partial B_1(x_0)$ and $t > 0$. We have

$$\begin{aligned} LW &= LU + \varepsilon\beta e^{\beta t} \\ &\geq b(x, t)U(x_0, t) + \varepsilon\beta e^{\beta t} \\ &= b(x, t)[W(x_0, t) - \varepsilon e^{\beta t}] + \varepsilon\beta e^{\beta t}. \end{aligned}$$

This gives

$$LW - b(x, t)W(x_0, t) \geq -\varepsilon b(x, t)e^{\beta t} + \varepsilon\beta e^{\beta t} = \varepsilon e^{\beta t}(\beta - b(x, t)).$$

By choosing $\beta > b(x, t)$ for $(x, t) \in \overline{B_1(x_0)} \times [0, \infty)$, we have

$$LW - b(x, t)W(x_0, t) > 0 \text{ in } B_1(x_0) \times (0, T).$$

If $W(x, t) \leq 0$ somewhere in $B_1(x_0) \times (0, T)$, then the set

$$\{t : W(x, t) \leq 0 \text{ for some } x \in B_1(x_0)\}$$

is nonempty. Let \tilde{t} be the infimum of the above set. Since $W(x, 0) > 0$ for $x \in \overline{B_1(x_0)}$, we have $0 < \tilde{t} < T$. At \tilde{t} , there exists some $x_1 \in B_1(x_0)$ such that $W(x_1, \tilde{t}) = 0$ and

$W_t(x_1, \tilde{t}) \leq 0$. On the other hand, W attains its local minimum at (x_1, \tilde{t}) . Therefore, $\Delta W(x_1, \tilde{t}) > 0$. Then, at $t = \tilde{t}$,

$$0 \geq W_t(x_1, \tilde{t}) > LW(x_1, \tilde{t}) - b(x_1, \tilde{t})W(x_0, \tilde{t}) > 0.$$

This gives a contradiction. Hence, $W(x, t) > 0$ in $B_1(x_0) \times (0, T)$. As $\varepsilon \rightarrow 0$, we have $U(x, t) \geq 0$ on $\overline{B_1(x_0)} \times [0, T)$. \square

By Lemma 2.1, 0 is a lower solution of the problem (1.1)-(1.2). On the other hand, u is bounded above by c . Thus, $0 \leq u < c$ on $\overline{B_1(x_0)} \times [0, T)$. Since u ceases to exist for $u \geq c$. It follows from Theorem 2.1 of [2] that the problem (1.1)-(1.2) has a unique classical solution $u \in C(\overline{B_1(x_0)} \times [0, T)) \cap C^{2+\alpha, 1+\alpha/2}(B_1(x_0) \times [0, T))$ for some $\alpha \in (0, 1)$ such that $0 \leq u < c$ on $\overline{B_1(x_0)} \times [0, T)$. From the hypothesis (H), f is differentiable. By the mean value theorem, there exists some positive constant k_1 (depending on a and f') such that

$$|af(u_1(x_0, t)) - af(u_2(x_0, t))| \leq k_1 |u_1(x_0, t) - u_2(x_0, t)|.$$

It follows from Theorem 8.9.2 of Pao [9, p. 436] that u either exists globally or only exists in a finite time.

To clarify the calculation in proving Lemma 3.3, we give the detail of the proof of the following lemma.

Lemma 2.2. *Let u be the solution of the problem (1.1)-(1.2). Then, $u_t \geq 0$ on $\overline{B_1(x_0)} \times [0, T)$ and $u_t > 0$ in $B_1(x_0) \times (0, T)$.*

Proof. For $h > 0$, let us consider the equation (1.1) at $t + h$. We have $Lu(x, t + h) = af(u(x_0, t + h))$ in $B_1(x_0) \times (0, T - h)$. Subtract equation (1.1) from this equation, and based on the mean value theorem, there exists some ζ_1 where ζ_1 is between $u(x_0, t + h)$ and $u(x_0, t)$ such that

$$Lu(x, t + h) - Lu(x, t) = af'(\zeta_1)[u(x_0, t + h) - u(x_0, t)] \text{ in } B_1(x_0) \times (0, T - h).$$

Since $u \geq 0$ on $\overline{B_1(x_0)} \times [0, T)$, we have $u(x, h) - u(x, 0) \geq 0$ for $x \in \overline{B_1(x_0)}$. From the boundary condition, $u(x, t + h) - u(x, t) = 0$ for $t > 0$ and $x \in \partial B_1(x_0)$. By Lemma 2.1, $u(x, t + h) \geq u(x, t)$ on $\overline{B_1(x_0)} \times [0, T - h)$. Thus, $(u(x, t + h) - u(x, t))/h \geq 0$ on $\overline{B_1(x_0)} \times [0, T - h)$. As $h \rightarrow 0^+$, $u_t \geq 0$ on $\overline{B_1(x_0)} \times [0, T)$.

From the equation (2.2) of Friedman [5, p. 61], the Hölder norm of $f'(u(x_0, t))$ with an exponent α is given by

$$|f'(u(x_0, t))|_\alpha = \sup_{[0, T)} |f'(u(x_0, t))| + \sup_{\substack{t \in [0, T) \\ \tilde{t} \in [0, T)}} \frac{|f'(u(x_0, t)) - f'(u(x_0, \tilde{t}))|}{[\sqrt{|t - \tilde{t}|}]^\alpha}.$$

By the mean value theorem, we obtain

$$|f'(u(x_0, t))|_\alpha \leq \sup_{[0, T)} |f'(u(x_0, t))| + \left(\sup_{B_1(x_0) \times [0, T)} |f''| \right) |u|_\alpha.$$

Thus, $|f'(u(x_0, t))|_\alpha$ is bounded when $0 \leq u < c$. From equation (3.2) of Friedman [5, p. 66], we obtain

$$|af'(u(x_0, t))u_t(x_0, t)|_\alpha \leq a|f'(u)|_\alpha |u_t|_\alpha.$$

Since $u \in C\left(\overline{B_1(x_0)} \times [0, T]\right) \cap C^{2+\alpha, 1+\alpha/2}(B_1(x_0) \times [0, T])$, we have

$$af'(u)u_t \in C^{\alpha, \alpha/2}(B_1(x_0) \times [0, T]).$$

It follows from Theorem 3.6 of Friedman [5, p. 65] that u_{tt} and $u_{x_i x_j t}$ exist for $i, j = 1, 2, \dots, n$. To show that u_t is positive, we differentiate equation (1.1) with respect to t to get

$$Lu_t = af'(u(x_0, t))u_t(x_0, t) \text{ in } B_1(x_0) \times (0, T).$$

For any $(x, t) \in B_1(x_0) \times [0, T]$, the integral representation form of u_t is given by

$$(2.1) \quad \begin{cases} u_t(x, t) = \int_{\overline{B_1(x_0)}} G(x, \xi, t) u_t(\xi, 0) d\xi \\ \quad + a \int_0^t \int_{\overline{B_1(x_0)}} G(x, \xi, t - \tau) f'(u(x_0, \tau)) u_t(x_0, \tau) d\xi d\tau, \end{cases}$$

where $G(x, \xi, t - \tau)$ is Green's function of the problem (1.1) subject to the homogeneous boundary condition. Since $f' > 0$, and $G(x, \xi, t - \tau) > 0$ in the set $\{(x, \xi, t) : x \text{ and } \xi \text{ are in } B_1(x_0), \text{ and } t > \tau \geq 0\}$, it follows from equation (2.1) that $u_t > 0$ in the domain $B_1(x_0) \times (0, T)$. \square

Without loss of generality, let us assume that x_0 is the origin 0. By the symmetry of $B_1(0)$, the polar form of the problem (1.1)-(1.2) is given by

$$(2.2) \quad \begin{cases} u_t(r, t) - u_{rr}(r, t) - \frac{n-1}{r}u_r(r, t) = af(u(0, t)) \text{ in } (0, 1) \times (0, T), \\ u(r, 0) = 0 \text{ for } r \in [0, 1], u_r(0, t) = 0 \text{ and } u(1, t) = 0 \text{ for } t \in (0, T). \end{cases}$$

Lemma 2.3. *The solution $u(r, t)$ of the problem (2.2) attains its maximum at $r = 0$ for $t \in (0, T)$.*

Proof. The solution of the problem (2.2) is radial symmetric with respect to $r = 0$. To show that $u_r < 0$ for $r \in (0, 1]$, let $H(r, t) = u_r(r, t)$. Differentiating the first equation of the problem (2.2) with respect to r , we have

$$H_t - H_{rr} - \frac{n-1}{r}H_r + \frac{n-1}{r^2}H = 0 \text{ for } (r, t) \in (0, 1) \times (0, T).$$

At $t = 0$, $H(r, 0) = 0$ for $r \in [0, 1]$. By Lemma 2.2, $u_t > 0$ in $B_1(0) \times (0, T)$. By Hopf's Lemma, $H(1, t) < 0$ for $t \in (0, T)$. Also, $H(0, t) = u_r(0, t) = 0$ for $t \in [0, T]$. By the maximum principle [9, p. 54], $H < 0$ for $(r, t) \in (0, 1] \times (0, T)$. Therefore, $u(0, t) \geq u(r, t)$ for $(r, t) \in [0, 1] \times (0, T)$. \square

3. SINGLE POINT QUENCHING OF u AND BLOW-UP OF u_t IN THE WHOLE DOMAIN

Let $\phi(x)$ be the eigenfunction corresponding to the first eigenvalue $\lambda_1 (> 0)$ of the Sturm-Liouville problem:

$$\begin{aligned}\Delta Z + \lambda Z &= 0 \text{ in } B_1(0), \\ Z(x) &= 0 \text{ for } x \in \partial B_1(0).\end{aligned}$$

By Theorem 3.1.2 of Pao [9, p. 97], $\phi(x) > 0$ in $B_1(0)$. The following lemma shows that u quenches in a finite time when a is sufficiently large, and give an upper bound of the quenching time.

Lemma 3.1. *If $af'(0) > \lambda_1$, then there exists some finite time \hat{T} such that u quenches as $t \rightarrow \hat{T}^-$ and*

$$(3.1) \quad \frac{1}{af'(0) - \lambda_1} \ln \left[\frac{c(af'(0) - \lambda_1) + af(0)}{af(0)} \right] \geq \hat{T}.$$

Proof. From Lemma 2.3, $u(0, t) \geq u(x, t)$ for $(x, t) \in \overline{B_1(0)} \times [0, T)$. Let $v(x, t)$ be the solution to the following initial-boundary value problem:

$$(3.2) \quad \begin{cases} Lv(x, t) = af(v(x, t)) \text{ for } (x, t) \in B_1(0) \times (0, T), \\ v(x, 0) = 0 \text{ for } x \in \overline{B_1(0)}, v(x, t) = 0 \text{ for } x \in \partial B_1(0) \text{ and } t > 0. \end{cases}$$

By the maximum principle, $v(x, t) > 0$ in $B_1(0) \times (0, T)$. Since $f' > 0$, we have $u \geq v$ on $\overline{B_1(0)} \times [0, T)$. We would like to prove that v quenches in some finite time when a is sufficiently large. By the hypothesis (H) and Taylor's Theorem, $f(v) = f(0) + f'(0)v + f''(\zeta_2)v^2/2$ for some ζ_2 between 0 and v . Since $f'' > 0$, we have $f(v) \geq f(0) + f'(0)v$. From the first equation of (3.2),

$$Lv \geq af(0) + af'(0)v.$$

Multiplying both sides by $\phi(x)$, integrating this inequality over $\overline{B_1(0)}$, and using the boundary conditions of v and ϕ , we obtain

$$\left(\int_{B_1(0)} v\phi dx \right)_t \geq af(0) \int_{B_1(0)} \phi dx + (af'(0) - \lambda_1) \int_{B_1(0)} v\phi dx.$$

Solving the above inequality on $[0, t]$, we have

$$\int_{B_1(0)} v\phi dx \geq \frac{af(0) \int_{B_1(0)} \phi dx [e^{(af'(0) - \lambda_1)t} - 1]}{af'(0) - \lambda_1}.$$

Since $af'(0) > \lambda_1$, the right side is an increasing function of t . Hence, $v(x, t)$ reaches c in some finite time \tilde{T} for some x in $B_1(0)$. We note that \tilde{T} satisfies

$$c \int_{B_1(0)} \phi dx \geq \frac{af(0) \int_{B_1(0)} \phi dx [e^{(af'(0) - \lambda_1)\tilde{T}} - 1]}{af'(0) - \lambda_1}.$$

Solving the above inequality for \tilde{T} , we obtain an upper bound of the quenching time for $v(x, t)$ stated in the inequality (3.1). Since $u \geq v$, the quenching time for u is $\hat{T} \leq \tilde{T}$. The lemma is proved. \square

From the result of Lemma 2.3, $u_r < 0$ for $(r, t) \in (0, 1] \times (0, T)$. This implies if u quenches, then $x = 0$ is the quenching point. Let T be the supremum of the time for which the problem (1.1)-(1.2) has a unique solution $u \in C\left(\overline{B_1(0)} \times [0, T)\right) \cap C^{2+\alpha, 1+\alpha/2}(B_1(0) \times [0, T))$. The following result shows that $u(0, t)$ quenches at T if T is finite.

Theorem 3.2. *If $T < \infty$, then $u(0, t)$ quenches at T .*

Proof. Suppose that u does not quench at $x = 0$ when $t = T$. There exists a positive constant k_2 such that $u(0, t) \leq k_2 < c$ for $t \in [0, T]$. This shows that $f(u(0, t)) < Q$ for some positive constant Q and $t \in [0, T]$. Then, by Theorem 4.2.1 of Ladde et al. [7, p. 139], $u \in C\left(\overline{B_1(0)} \times [0, T]\right) \cap C^{2+\alpha, 1+\alpha/2}(B_1(0) \times [0, T])$. This implies that there exists a positive constant k_3 such that $u(x, t) \leq k_3 < c$ for $(x, t) \in \overline{B_1(0)} \times [0, T]$. To arrive at a contradiction, we need to show that u can be continued into a larger time interval $[0, T + t_1)$ for some positive t_1 . This can be achieved by extending the upper bound. Let us construct an upper solution $\psi(x, t) = k_3 m(t)$, where $m(t)$ is a positive function and satisfies

$$\frac{d}{dt}m(t) = \frac{a}{k_3}f(k_3 m(t)) \text{ for } t \geq T, \quad m(T) = 1.$$

Because $f(u) \in C^2([0, c])$, the solution $m(t)$ exists. Since $f > 0$, $m(t)$ is an increasing function of t . Let t_1 be a positive constant determined by $k_3 m(T + t_1) = k_4 < c$ for some positive constant k_4 greater than k_3 . By construction, $\psi(x, t)$ is the solution to the following problem:

$$\begin{aligned} L\psi &= af(\psi) \text{ in } B_1(0) \times [T, T + t_1), \\ \psi(x, T) &= k_3 \geq u(x, T) \text{ on } \overline{B_1(0)}, \\ \psi(x, t) &= k_3 m(t) > 0 \text{ for } x \in \partial B_1(0) \text{ and } t \in (T, T + t_1). \end{aligned}$$

By Lemma 2.1, $\psi(x, t) \geq u(x, t)$ on $\overline{B_1(0)} \times [T, T + t_1)$. Therefore, we find a solution u to the problem (1.1)-(1.2) on $\overline{B_1(0)} \times [T, T + t_1)$. This contradicts the definition of T . Hence, $u(0, t)$ quenches at T . \square

Let $z = u_t$. Differentiating the differential equation in (2.2) with respect to t , we have

$$(3.3) \quad z_t(r, t) - z_{rr}(r, t) - \frac{(n-1)}{r}z_r(r, t) = af'(u(0, t))z(0, t).$$

Using equation (3.3), we have

$$(3.4) \quad \begin{cases} z_t - z_{rr} - \frac{(n-1)}{r} z_r = af'(u(0, t)) z(0, t) & \text{for } (r, t) \in (0, 1) \times (0, T), \\ z(r, 0) \geq 0 \text{ for } r \in [0, 1], z(0, t) > 0 \text{ and } z(1, t) = 0 & \text{for } t \in (0, T). \end{cases}$$

The following result shows that u_t attains its maximum at $r = 0$.

Lemma 3.3. $u_t(r, t) < u_t(0, t)$ for $(r, t) \in (0, 1] \times (0, T)$.

Proof. Using a similar calculation as in the proof of Lemma 2.2, we obtain u_{ttr} , u_{rt} , u_{rrt} , and u_{rrrt} . By differentiating the first equation of the problem (3.4) with respect to r , we obtain

$$z_{tr} - z_{rrr} - \frac{(n-1)}{r} z_{rr} + \frac{(n-1)}{r^2} z_r = 0.$$

For $r \in (0, 1)$, $u_{rt}(r, 0)$ is given by

$$u_{rt}(r, 0) = \lim_{h \rightarrow 0} \frac{u_r(r, h) - u_r(r, 0)}{h}.$$

Using $u_r(r, 0) = 0$ and Lemma 2.3, we have $u_{rt}(r, 0) \leq 0$. Thus, $z_r(r, 0) \leq 0$ for $r \in (0, 1)$. By Lemma 2.2,

$$\frac{\partial z(1, 0)}{\partial r} = \lim_{h \rightarrow 0} \frac{z(1, 0) - z(1-h, 0)}{h} \leq 0.$$

By Hopf's lemma, $\partial z(1, t)/\partial r < 0$ for $t > 0$. By the symmetry of $B_1(0)$ with respect to $r = 0$, $\partial z(0, t)/\partial r = 0$ for $t \geq 0$. Let $V = z_r$. Then, V satisfies the following initial-boundary value problem:

$$V_t - V_{rr} - \frac{(n-1)}{r} V_r + \frac{(n-1)}{r^2} V = 0 \text{ for } (r, t) \in (0, 1) \times (0, T),$$

$$V(r, 0) \leq 0 \text{ for } r \in [0, 1], \text{ and } V(0, t) = 0 \text{ and } V(1, t) < 0 \text{ for } t \in (0, T).$$

By the maximum principle, $V(r, t) < 0$ for $(0, 1] \times (0, T)$. Integrating $V(r, t) < 0$ over $(0, r)$, we obtain $z(r, t) < z(0, t)$, that is, $u_t(r, t) < u_t(0, t)$ for $(r, t) \in (0, 1] \times (0, T)$. \square

We modify Theorem 5 of Chan and Liu [3] to prove the following theorem.

Theorem 3.4. $u(r, t)$ quenches only at $r = 0$.

Proof. Let $I(r, t) = r^{n-1} u_r(r, t)$. From the result of Lemma 2.3, $I(r, t) < 0$ for $(r, t) \in (0, 1] \times (0, T)$. From the first equation of the problem (2.2), $I(r, t)$ satisfies the following differential equation,

$$u_t - \frac{1}{r^{n-1}} I_r = af(u(0, t)).$$

Differentiating the above equation with respect to r , we get

$$u_{rt} - \frac{1}{r^{n-1}} I_{rr} + \frac{(n-1)}{r^n} I_r = 0.$$

By rewriting the above equation in terms of I , we have

$$\frac{I_t}{r^{n-1}} - \frac{1}{r^{n-1}} I_{rr} + \frac{(n-1)}{r^n} I_r = 0.$$

That is,

$$(3.5) \quad I_t - I_{rr} + \frac{(n-1)}{r} I_r = 0.$$

From Lemma 2.3, there exists a number $\delta \in (0, 1)$ such that $u_r < 0$ for $(r, t) \in (1 - \delta, 1] \times (0, T)$. Thus, there exist positive real numbers $t_2 (< T)$ and μ such that $I(r, t_2) < -\mu$ for $r \in (1 - \delta, 1)$. Let r_1 be an arbitrary positive real number such that $1 - \delta < r_1 < 1$, κ be a positive number, and $R(r, t) = I(r, t) + \kappa(1 - r)$. At $t = t_2$, let us choose κ such that $R(r, t_2) < -\mu + \kappa(1 - r) < 0$ for $r \in [r_1, 1)$. By Lemma 3.3, $u_{rt}(r_1, t) < 0$ for $t \in [t_2, \omega)$ where $\omega \leq T$. That is, $u_r(r_1, t)$ is a decreasing function of t . Since $r_1^{n-1} u_r(r_1, t_2) < -\mu$, we then obtain $r_1^{n-1} \inf_{[t_2, T]} u_r(r_1, t) + \kappa(1 - r_1) < -\mu + \kappa(1 - r_1) < 0$. Clearly, $R(1, t) < 0$ for $t \in (0, T)$. Also, $R_r = I_r - \kappa$, $R_{rr} = I_{rr}$, and $R_t = I_t$. With these, we can rewrite equation (3.5) in terms of R as

$$R_t - R_{rr} + \frac{(n-1)}{r} R_r = \frac{-\kappa(n-1)}{r} < 0.$$

By the maximum principle, $R < 0$ for $(r, t) \in [r_1, 1] \times [t_2, T)$. That is, $r^{n-1} u_r + \kappa(1 - r) < 0$. Equivalently,

$$u_r < -\kappa \frac{(1-r)}{r^{n-1}}.$$

Integrating both sides from r_2 to r_3 where $r_1 < r_2 < r_3 < 1$, we have

$$u(r_3, t) - u(r_2, t) < -\kappa \int_{r_2}^{r_3} \frac{(1-r)}{r^{n-1}} dr.$$

Since $\int_{r_2}^{r_3} (1-r)/r^{n-1} dr > 0$,

$$\begin{aligned} u(r_3, t) &< u(r_2, t) - \kappa \int_{r_2}^{r_3} \frac{(1-r)}{r^{n-1}} dr \\ &< c - \kappa \int_{r_2}^{r_3} \frac{(1-r)}{r^{n-1}} dr, \end{aligned}$$

which shows that $u(r_3, t) < c$ for $t_2 \leq t < T$. Thus, u does not quench at r_3 for any $r_3 \in (r_1, 1)$. Because $\delta \in (0, 1)$ and r_1 is an arbitrary real number such that $1 - \delta < r_1 < 1$, we have $r_1 \in (0, 1)$. Therefore, for any $r_3 \in (r_1, 1) \subseteq (0, 1)$, r_3 is not a quenching point. Hence, the solution $u(r, t)$ quenches only at $r = 0$. \square

Let $\varphi_0(x) \in C(\overline{B_1(0)}) \cap C^2(B_1(0))$ such that $\Delta \varphi_0(x) < 0$, $\varphi_0(x) > 0$ in $B_1(0)$, and $\varphi_0(x) = 0$ for $x \in \partial B_1(0)$ and $\max_{x \in \overline{B_1(0)}} \varphi_0(x) \leq 1$. Let $\varphi(x, t)$ be the solution to the following first initial-boundary value problem:

$$\begin{aligned} Lw &= 0 \text{ in } B_1(0) \times (0, \infty) \\ w(x, 0) &= \varphi_0(x) \text{ on } \overline{B_1(0)} \\ w(x, t) &= 0 \text{ for } (x, t) \in \partial B_1(0) \times (0, \infty). \end{aligned}$$

By the maximum principle, $\varphi(x, t) > 0$ in $B_1(0) \times (0, \infty)$ and is bounded above by $\varphi_0(x)$. Further,

$$\max_{(x,t) \in \overline{B_1(0)} \times [0, \infty)} \varphi(x, t) \leq 1.$$

By Lemma 2.2, $u_t(x, t) > 0$ in $B_1(0) \times (0, T)$. From equation (3.3), we have

$$z_t(r, t) - z_{rr}(r, t) - \frac{(n-1)}{r} z_r(r, t) = af'(u(0, t)) z(0, t).$$

By Hopf's lemma, $\partial z / \partial r < 0$ at $r = 1$ for $t > 0$. Then for any positive $\nu (< T)$, there exists a positive number σ depending on ν such that $\partial z / \partial r \leq -\sigma$ at $t = \nu$ for $r \in (1 - \gamma, 1)$, where γ is some positive number less than 1. Integrating $\partial z / \partial r \leq -\sigma$ over $(r, 1)$ at $t = \nu$, we obtain

$$\int_{z(r, \nu)}^{z(1, \nu)} dz \leq \int_r^1 -\sigma dr,$$

which gives

$$z(1, \nu) - z(r, \nu) \leq -\sigma(1 - r).$$

From $z(1, \nu) = 0$, we obtain $\sigma(1 - r) \leq z(r, \nu)$. Let $\rho = \min_{x \in \overline{B_{1-\gamma}(0)}} u_t(x, \nu)$. Then, ρ is a positive constant. Let us choose $\eta_1 \in (0, 1)$ such that

$$\rho \geq \eta_1 \varphi(x, \nu) f(u(0, \nu)) \text{ for } x \in \overline{B_{1-\gamma}(0)}.$$

For $x \in B_1(0) \setminus \overline{B_{1-\gamma}(0)}$, there exists $\eta_2 \in (0, 1)$ such that

$$\sigma(1 - r) \geq \eta_2 \varphi(r, \nu) f(u(0, \nu)) \text{ for } r \in (1 - \gamma, 1),$$

where $r = \|x\|$. From this, we get

$$z(x, \nu) = u_t(x, \nu) \geq \eta_2 \varphi(x, \nu) f(u(0, \nu)) \text{ for } x \in B_1(0) \setminus \overline{B_{1-\gamma}(0)}.$$

Let $\eta = \min\{\eta_1, \eta_2\}$. We obtain

$$u_t(x, \nu) \geq \eta \varphi(x, \nu) f(u(0, \nu)) \text{ for } x \in B_1(0).$$

For $x \in \partial B_1(0)$, the above inequality becomes an equality, and the left and right sides equal. Therefore,

$$(3.6) \quad u_t(x, \nu) \geq \eta \varphi(x, \nu) f(u(0, \nu)) \text{ for } x \in \overline{B_1(0)}.$$

Let $J(x, t) = u_t(x, t) - \eta \varphi(x, t) f(u(0, t))$. We modify the proof of Lemma 3.4 of [4] to obtain the following result.

Lemma 3.5. *For a fixed positive number $\nu (< T)$, $J(x, t) \geq 0$ on $\overline{B_1(0)} \times [\nu, T)$.*

Proof. We have

$$J_t = u_{tt} - \eta \varphi_t f(u(0, t)) - \eta \varphi f'(u(0, t)) u_t(0, t),$$

$$\Delta J = \Delta u_t - \eta f(u(0, t)) \Delta \varphi.$$

Then,

$$\begin{aligned} LJ &= Lu_t - \eta f(u(0, t)) L\varphi - \eta \varphi f'(u(0, t)) u_t(0, t) \\ &= f'(u(0, t)) u_t(0, t) (1 - \eta \varphi). \end{aligned}$$

By $f' > 0$, $\varphi \leq 1$ and $\eta < 1$, and $u_t(0, t) > 0$ for $t \in (0, T)$, we obtain $LJ \geq 0$ in $B_1(0) \times (0, T)$. By equation (3.6), $J(x, \nu) \geq 0$ on $\overline{B_1(0)}$. By the boundary conditions, $J(x, t) = 0$ on $\partial B_1(0) \times [\nu, T)$. By the maximum principle, $J(x, t) \geq 0$ on $\overline{B_1(0)} \times [\nu, T)$. \square

Based on Lemma 3.5, we obtain the following corollary.

Corollary 3.6. *If $u(0, t)$ quenches at $t = T$, then $u_t(x, t) \rightarrow \infty$ as $t \rightarrow T^-$ for all $x \in B_1(0)$.*

Proof. From the result of Lemma 3.5, $u_t(x, t) \geq \eta \varphi(x, t) f(u(0, t))$ in $B_1(0) \times [\nu, T)$. If $u(0, t) \rightarrow c^-$ as $t \rightarrow T^-$, then $f(u(0, t)) \rightarrow \infty$. Because $0 < \varphi(x, t) \leq 1$ in $B_1(0) \times (0, \infty)$, we have $u_t(x, t) \rightarrow \infty$ as $t \rightarrow T^-$ for all $x \in B_1(0)$. \square

4. CONCLUSION

In this paper, we prove that $u_t(x, t)$ blows up in the whole domain of an n -dimensional ball if $u(x_0, t)$ quenches in some finite time. The technique used is through determining a lower bound of u_t . This lower bound tends to infinity when $u(x_0, t)$ quenches. A necessary condition for quenching of u is also given. Also, we show that $u(x, t)$ quenches at x_0 only in some finite time.

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REFERENCES

- [1] K. Bimpong-Bota, P. Ortoleva, and J. Ross, Far-from-equilibrium phenomena at local sites of reaction, *J. Chem. Phys.*, 60:3124–3133, 1974.
- [2] J. M. Chadam, A. Peirce, and H. M. Yin, The blowup property of solutions to some diffusion equations with localized nonlinear reactions, *J. Math. Anal. Appl.*, 169:313–328, 1992.
- [3] W. Y. Chan and H. T. Liu, Finding the critical domain of multi-dimensional quenching problems with Neumann boundary conditions, *Neural Parallel Sci. Comput.*, 25:19–28, 2017.
- [4] C. Chang, Y. Hsu, and H. T. Liu, Quenching behavior of parabolic problems with localized reaction term, *Mathematics and Statistics*, 2:48–53, 2014.
- [5] A. Friedman, *Partial Differential Equations of Parabolic Type*, Prentice-Hall, Englewood Cliffs, NJ, 1964, pp. 61, and 65–66.
- [6] H. Kawarada, On solutions of initial-boundary problem for $u_t = u_{xx} + 1/(1 - u)$, *Publ. Res. Inst. Math. Sci., Kyoto Univ.*, 10:729–736, 1975.
- [7] G. S. Ladde, V. Lakshmikantham, and A. S. Vatsala, *Monotone Iterative Techniques for Nonlinear Differential Equations*, Pitman, 1985, p. 139.
- [8] P. Ortoleva and J. Ross, Local structures in chemical reactions with heterogeneous catalysis, *J. Chem. Phys.*, 56:4397–4400, 1972.
- [9] C. V. Pao, *Nonlinear Parabolic and Elliptic Equations*, Plenum Press, New York, NY, 1992, pp. 54, 97, and 436.